

Kelvin-Helmholtz instability and magnetic reconnection at the Earth's magnetopause : 3D simulation based on satellite data

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Outline

Introduction

- The magnetospheric environment
- KH as a large-scale (MHD) flute mode
- Type I Vortex Induced Reconnection
- Mid-latitude reconnection
- MMS observations & motivations

Setting-up simulations

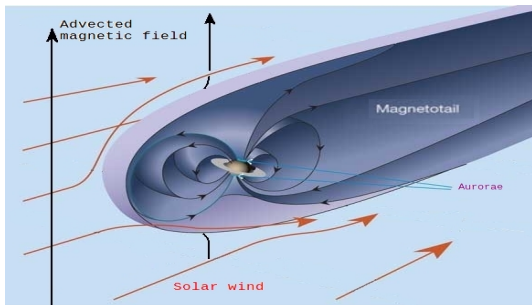
- A model for the magnetospheric flank
- High-latitude stabilization & magnetic rotation

3D “realistic” simulation results

- KH and current sheet dynamics
- Reconnection dynamics and its latitude distribution
- In-situ* and *remote* events: simulation Vs observations

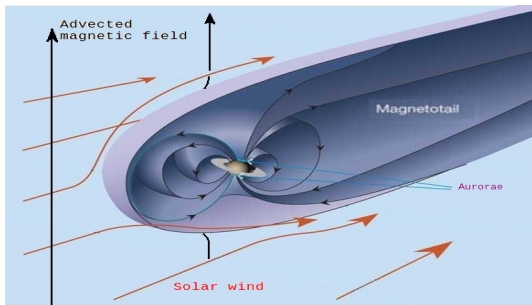
Conclusions

Magnetospheric environment: Northward “quiet” periods



- ▶ Unexpected **efficient transport** between the solar wind and the magnetosphere
- ▶ **Inferred** $D_{eff} \simeq 10^9 m^2/s$
- ▶ **Cross-field diffusivity** (collisional or anomalous) **too small**

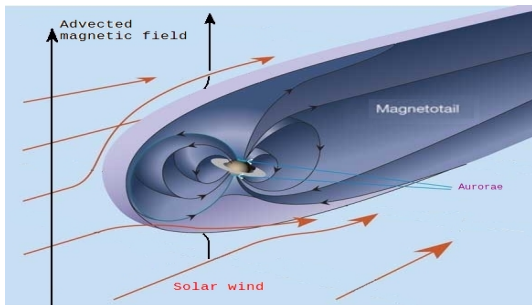
Magnetospheric environment: Northward “quiet” periods



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 - ▶ **Inferred** $D_{eff} \simeq 10^9 m^2/s$
 - ▶ **Cross-field diffusivity** (collisional or anomalous) **too small**
 - ▶ Different mechanisms have been proposed →
- 1) **“Double lobe reconnection”** can generate a Low Latitude Boundary Layer, but it is **not sufficient**.^a
 - 2) **Kinetic Alfvén Waves** at the magnetopause can strongly enhance cross field diffusion, but **not often observed** and sometimes excluded^b

^aWing 06, Taylor 08 & Hasegawa 09 ^b Nakai 01 ^c Kavosi 15

Magnetospheric environment: Northward “quiet” periods

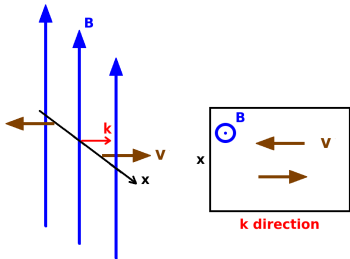


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 - 2) **Kinetic Alfvén Waves** at the magnetopause can strongly enhance cross field diffusion, but **not often observed** and sometimes excluded^b
 - 3) **Kelvin-Helmholtz instability**:
 - ▶ **“Robust”** phenomenon^c
 - ▶ A good **driver** for a **rich dynamics**

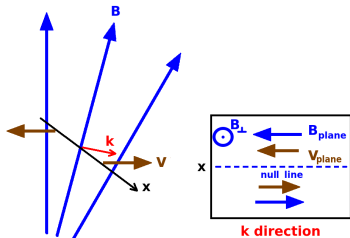
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KHI as a large-scale (MHD) flute mode

- ▶ KHI is **the instability** of a **sheared velocity configuration**
- ▶ Velocity shear half-width $a \sim$ several $d_i \rightarrow$ KHI develops as a large-scale (\sim tenths of d_i), **nearly magnetohydrodynamic mode**
- ▶ $\vec{k} \cdot \Delta \vec{V}$ provides the energy source, while $\vec{k} \cdot \vec{B}$ is the sink (**magnetic tension is stabilizing**)
- ▶ For **magnetospheric parameters**, KHI develops as a **flute mode**, with $\vec{k} \cdot \vec{B} \simeq 0$

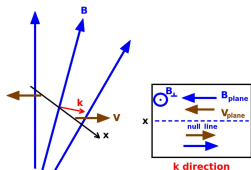


Without magnetic rotation
 $\vec{k} \cdot \vec{B} = 0$ everywhere.



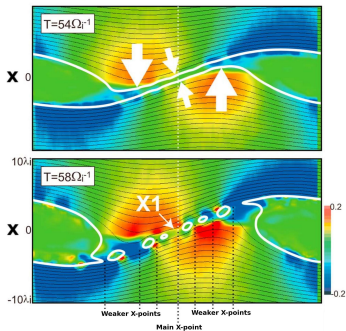
With magnetic rotation
 $\vec{k} \cdot \vec{B} = 0$ where $\partial_x V$ is max (**at magnetopause**)

Current pinching: Vortex Induced Reconnection



- ▶ $\partial_x V_{plane} > \partial_x V_{A,plane} \Rightarrow$ **KHI pinches the original current sheet and force reconnection to occur**
 \rightarrow **Type I Vortex Induced Reconnection.**

- ▶ **VIR creates field lines crossing the original frontier** between the magnetospheric and SW plasmas.

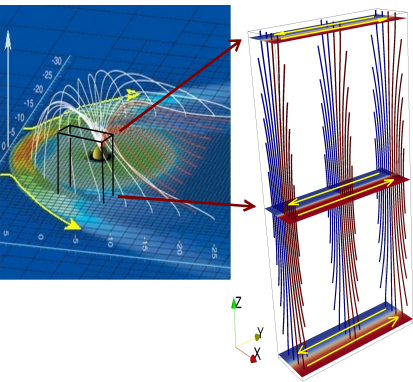


- ▶ **3D kinetic simulations show the streaming of particles along reconnected lines** \rightarrow formation of a **mixing layer^a**.

- ▶ Effective $D_{eff} = O(10^9 m^2/s)$ or **even higher^b**.

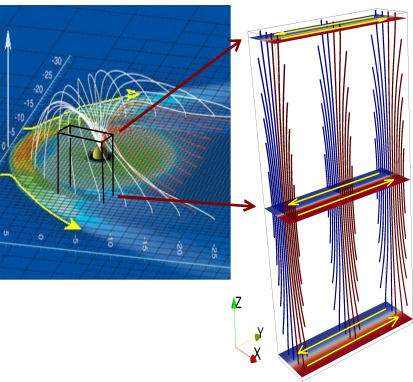
^aNakamura JGR 13; ^bNakamura Nat 17
 Bottom figures from Nakamura JGR 11

Current creation by high-latitude stabilization



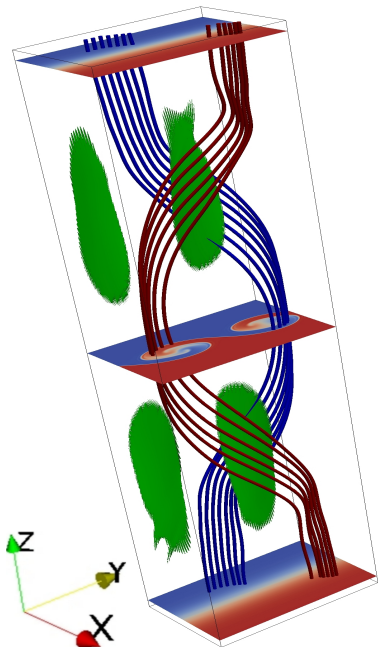
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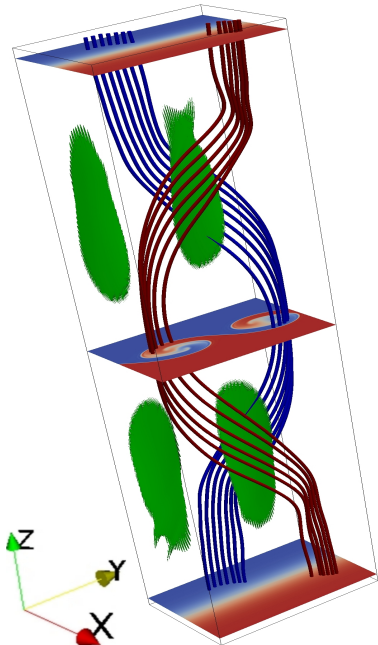
- ▶ **Vortices create the conditions** for reconnection **even if**, at low-latitude, there is **no magnetic rotation** (no initial current)
- ▶ The **equatorial** region is **unstable** since:
 - **velocity shear exists**
 - **reduced magnetic tension**
- ▶ **High latitudes**
 - complex configuration
 - **total stabilization**

Current sheets at mid-latitudes (resistive Hall-MHD sim.)



- ▶ **Low-latitude** region → **vortices**
- ▶ **High-latitude** regions → **stable**

Current sheets at mid-latitudes (resistive Hall-MHD sim.)



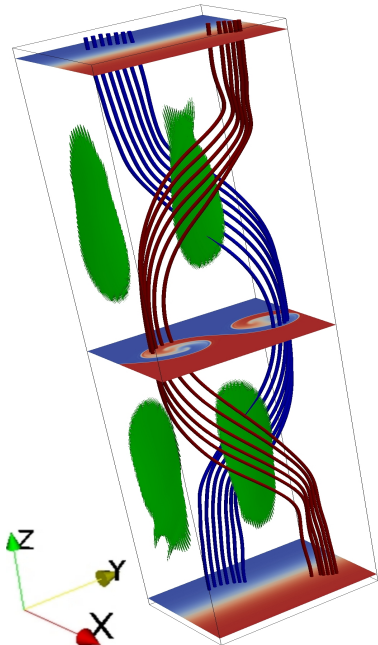
- ▶ **Low-latitude** region → **vortices**
- ▶ **High-latitude** regions → **stable**
- ▶ **Differential advection** for field lines
 - at v_{Solar_Wind} or $v_{Magnetosphere}$ at high latitudes
 - at $v_{phase} \simeq (v_{SW} + v_{Msph})/2$ at low latitude

⇒ **Arched solar wind** & **magnetospheric** field lines

⇒ **Mid-latitude current sheets**

→ **Favorable** conditions for **reconnection** to occur

Current sheets at mid-latitudes (resistive Hall-MHD sim.)



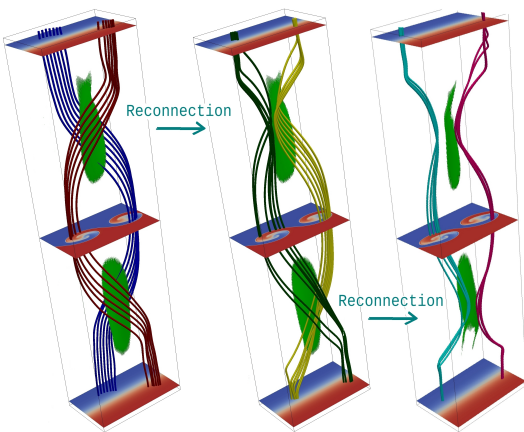
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Mid-latitude reconnection



- ▶ **Reconnection** occurs in **both hemispheres**
- ⇒ Creates **double reconnected lines**
 - ▶ They connect
N pole → **red arm** → **N pole**
- Flux tubes “closed” on the Earth populated by solar-wind particles**
(“Opened” flux tubes too.)
- ⇒ **Solar wind particles enter the magnetosphere**
- ▶ Effective^a $D_{eff} \simeq 10^9 m^2/s$
- ▶ Specific **entropy increases**^b

^aFaganello EPL 2012; ^b Johnson JGR 09

Motivations

- ▶ Signatures of **VIR** has been **measured** by **satellites**, close to the **equatorial plane**¹.
- ▶ **Signatures of Mid-latitude reconnection** have been found too²
- ▶ On **September 8th 2015**, **MMS** detected signatures of **both VIR** and **Mid-latitude reconnection**³
- ▶ It **observed an asymmetric distribution of "remote" reconnection events**, with **more events southern of the satellites**³
- ▶ We want to understand the **role** of the **two mechanism**, their **competition/cooperation** in determining the **magnetic evolution** and the **transport** properties of the **SW/magnetosphere frontier**

A model for the magnetospheric flank: 2D equilibrium with translation symmetry along the flow direction (y-direction)

- ▶ Ideal **MHD** equations & **adiabatic** closure
- ▶ 2D (x,z) equilibrium configuration:

$$\vec{B} = B_y \hat{e}_y + \nabla \psi \times \hat{e}_y, \quad \psi = \psi(x, z)$$

- ▶ Actually a distorted 1D equilibrium configuration:

$$B_y = B_y(\psi), \quad V = V_y(\psi) \hat{e}_y, \quad \rho = \rho(\psi), \quad p = p(\psi)$$

- ▶ **Simplified Grad-Shafranov** equation:

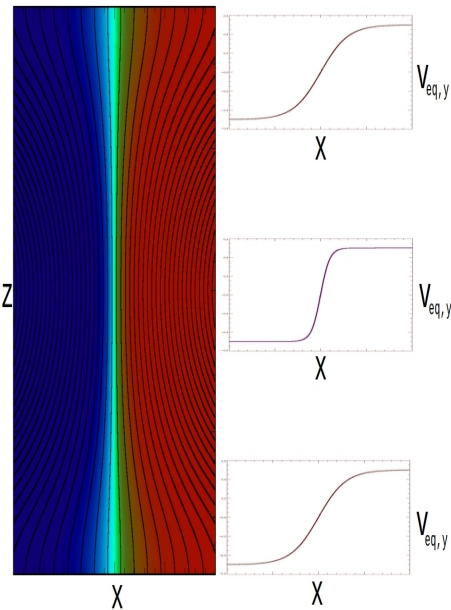
$$\nabla^2 \psi = -4\pi \frac{d\Pi}{d\psi}, \quad \Pi = \Pi(\psi) = p + \frac{B_y^2}{8\pi}$$

- ▶ Π uniformity \rightarrow **Laplace equation:** $\nabla^2 \psi = 0$

- ▶ A **suitable solution:**

$$\psi(x, z) = 1/2 [(1 + A)x + (1 - A)L_z/2\pi \sinh(2\pi x/L_z) \cos(2\pi z/L_z)]$$

High-latitude stabilization



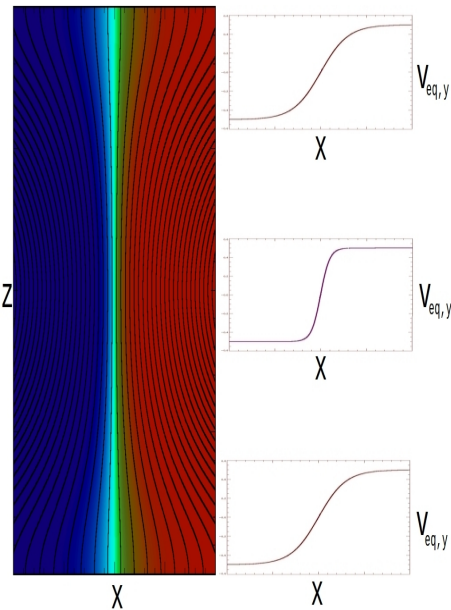
- ▶ **Hourglass-like ψ -isosurfaces**
(provide the dominant z-component of the magnetic field)

- ▶ $V_{eq,y} = \Delta V_{eq}/2 \tanh(\psi/a)$

→ **Stronger gradient at the equators**

^aFaganello *PPCF* 2012

High-latitude stabilization



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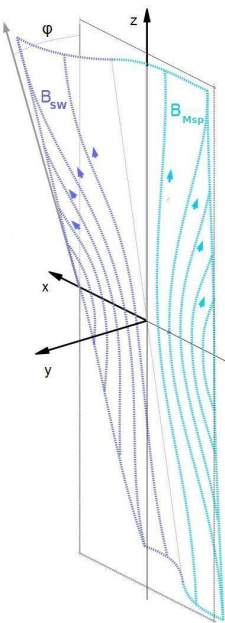
→ **Stronger gradient at the equators**

▶ $\gamma_{KH} \propto$ gradient

⇒ **High-Latitude stabilization^a**

^aFaganello *PPCF* 2012

High-latitude stabilization & magnetic rotation

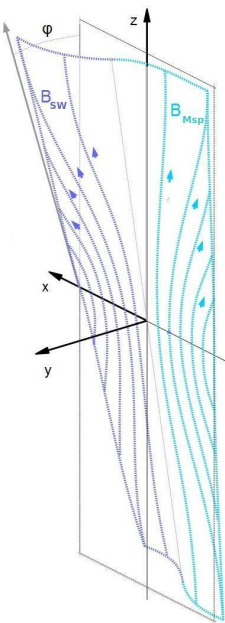


▶ $B_{eq,y} = B_{\parallel flow} (1 + \tanh(\psi/a))$

→ **Magnetic rotation**

^aFadanelli JGR 18

High-latitude stabilization & magnetic rotation



▶ $B_{eq,y} = B_{\parallel flow} (1 + \tanh(\psi/a))$

→ **Magnetic rotation**

▶ **Magnetic rotation & High-latitude stabilization**

⇒ **VIR and Mid-latitude reconnection, both at play.**

▶ $B_{\parallel flow}$ **breaks the N-S symmetry is broken**, for KH vortices and reconnection too^a.

^aFadanelli JGR 18

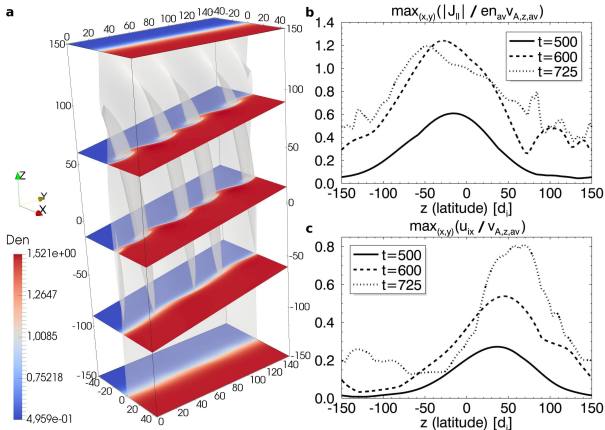
Data-based parameters for Hall-MHD simulation

- ▶ $\Delta V_{eq} \simeq 350 \text{ Km/s}$
- ▶ $B_z \simeq 67 \text{ nT}$
- ▶ $B_{flow} \simeq -20 \text{ nT}$
- ▶ $n \sim 5.7 \text{ cm}^{-3} \longrightarrow 20.1 \text{ cm}^{-3}$
- ▶ $T_{th} \sim 2800 \text{ eV} \longrightarrow 200 \text{ eV}$
- ▶ Shear half-width $a \simeq 900 \text{ Km}$ so that $\lambda_{KH} \sim 12000 \text{ Km}$

- ▶ $L_y = 2\lambda_{KH}$ (two vortices in the box) ; $n_y = 512$
- ▶ $L_z = 8\lambda_{KH} \sim 10^5 \text{ Km}$ i.e. $\pm 45^\circ$; $n_z = 512$
- ▶ $L_x = 2.4\lambda_{KH}$; $n_x = 900$ $\tau_{res,eq} = a^2/\eta \simeq 20 t_{max,sim}$

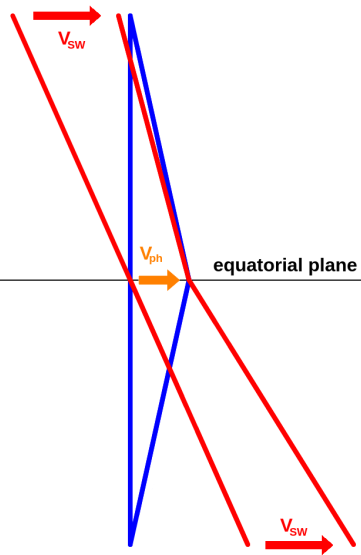
- ▶ Periodic boundary conditions (BC) along $y - z$ directions.
- ▶ BC along x -direction based on MHD characteristics:
 - Transparent boundaries
 - Sustains the 2D equilibrium

KH and current dynamics (resistive Hall-MHD simulation)



- ▶ **Oblique** \vec{k} , with $\vec{k} \cdot \vec{B}|_{magnetopause} \simeq 0$
- ▶ High-latitude stabilization
- ▶ **Asymmetric evolution**
- ▶ Folded magnetopause ($\psi = 0$ isosurface)

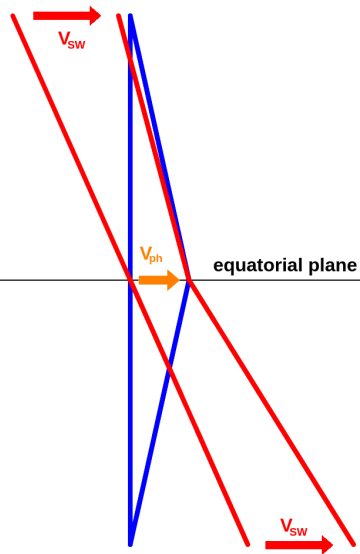
Differential advection & N-S asymmetry



- ▶ As soon as **field lines** are **caught into** the **vortices**, field lines are **advected differently** at **low/high-latitudes**
- ▶ As a consequence **magnetic rotation is enhanced** in **one hemisphere** (the southern one for $B_{\parallel flow} < 0$) while it is **lowered** in the **opposite one**.
- ▶ KH **vortices shift** where **rotation lowers** (smaller stabilization)

^aVernisse JGR 16, Vernisse JGR 20

Differential advection & N-S asymmetry

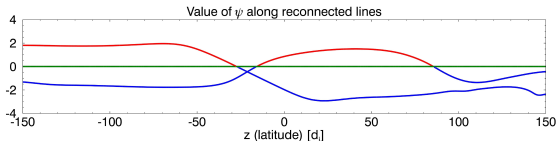


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- ▶ KH **vortices shift** where **rotation lowers** (smaller stabilization)
- ▶ Reconnection occurs as **both VIR** and **Mid-latitude one** but prefers the hemisphere where rotation (and J) rises.
- ▶ We **expect more reconnection** events in the **southern hemisphere**, as **observed by MMS** on 08/09/2015^a.

^aVernisse JGR 16, Vernisse JGR 20

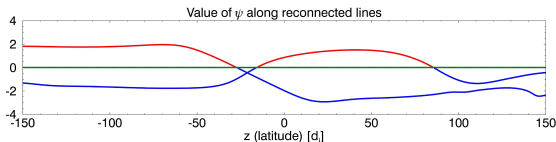
How to find reconnected lines in 3D, complex simulations ?

- ▶ ψ and \vec{B} are **advected independently**:
 - ψ follows an **ideal evolution**
 - \vec{B} the **true** resistive Hall-MHD one and **reconnects**.
- ▶ A jump of ψ along a line indicates reconnection.
- ▶ Current sheets, and reconnection, are at the magnetopause.
- ▶ Define a **line** as **“reconnected”** if $|\Delta\psi|_{\text{along line}} > a/2$ **across the magnetopause** ($\psi = 0$ isosurface).



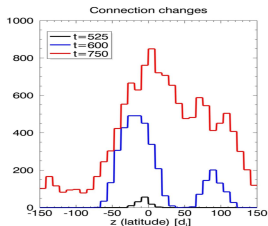
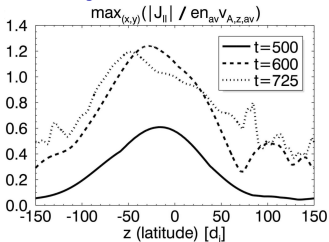
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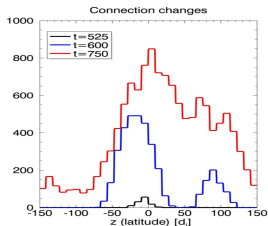
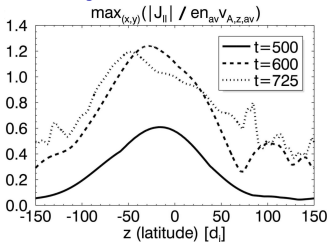
- ▶ 22500 lines integrated at each time → **statistical analysis**
- ▶ Define **reconnection event** as a **crossing** of the **magnetopause** by a reconnected line.

Reconnection dynamics and its latitude distribution



- ▶ **Early nonlinear evolution** ($t = 500\Omega_{ci}^{-1}$ and $600\Omega_{ci}^{-1}$):
 - The **original current sheet** is **compressed** around the **equators**
→ **VIR**
 - This **main current sheet** gradually **shifts southward** as well as **VIR** events
 - A **second current peak** arises in the **northern hemisphere**
→ **mid-latitude reconnection**.

Reconnection dynamics and its latitude distribution



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 - This **main current sheet** gradually **shifts southward** as well as **VIR** events
 - A **second current peak** arises in the **northern hemisphere**
→ **mid-latitude reconnection**.
- ▶ **Late nonlinear evolution** ($t = 725$): secondary small-scale KH vortices grows in between the two main current peaks. Reconnection is forced there and a wider distribution is formed.

Number of *in-situ* and *remote* events

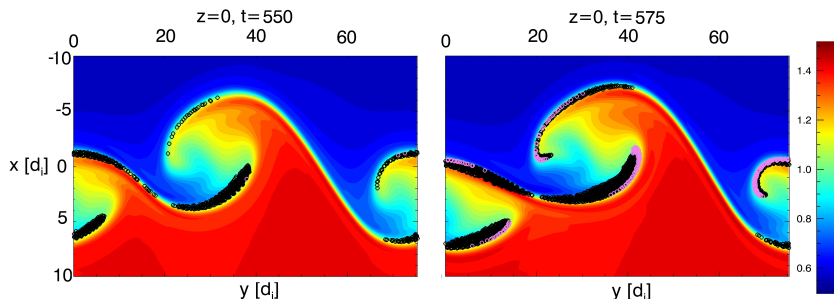
- ▶ Defined **with respect** to a virtual **satellite** at the **equators**.

Time interval	Southern h.	In-situ	Northern h.
501-525	35	81	2
526-550	377	201	7
551-575	677	321	327
576-600	730	268	590
601-625	544	145	686
626-650	694	247	953
651-675	488	356	1346
676-700	147	328	747
701-725	-	120	757

- ▶ **Early nonlinear** phase, as in the **MMS event**¹:
more remote events in the **southern hemisphere**
- ▶ We expect more *remote* events in the northern hemisphere if satellites would be further downstream along the flank

¹Vernisse JGR 20

Topology chart in the equatorial plane: satellite comparison



- ▶ **Black circles: crossing points of once-reconnected lines**
→ **expected *in-situ* signatures or *remote* signatures coming from the south.**
- ▶ **Purple circles: crossing points of double-reconnected lines**
→ **also *remote* signatures coming from the north.**
- ▶ In very **good agreement** with particle **distribution functions** observed by **MMS¹**.

¹Eriksson Front 2021

Conclusions

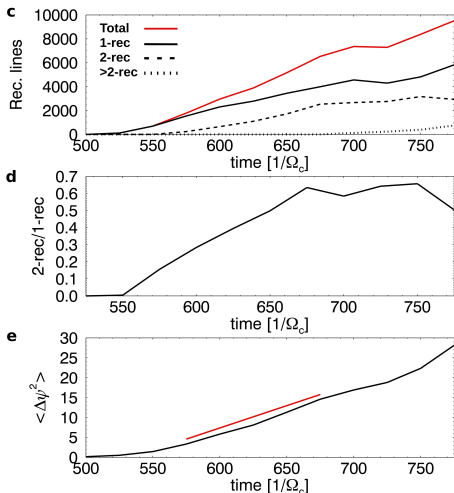
- ▶ **VIR** and **mid-latitude reconnection coexist** when an **initial magnetic rotation** is present.
- ▶ They lead to a **complex magnetic evolution** with an **asymmetric distribution of reconnection** events¹.
- ▶ The **predicted locations** for magnetic **reconnection** are in **good agreement** with reconnection **signatures** observed by **MMS**².
- ▶ **Even if** the system evolves **asymmetrically**, the number of **double-reconnected lines** reaches **40% of reconnected lines**¹ and **could explain** the **specific entropy increasing** that is observed **across the magnetopause**³.
- ▶ The **effective diffusion coefficient** associated to reconnection is **large enough** for explaining the **observed transport**¹

¹Sisti GRL 19, Faganello PPCF 22; ²Vernisse JGR 16, Eriksson Front 21; ³Johnson JGR 09

Thank you for your attention.

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Once/double reconnected lines and effective diffusion



- ▶ # of reconnected lines growth with time
 - ▶ # of **double-reconnected** lines reaches **70%** of # of **once-reconnected**.
 - ▶ ψ **measures** the **distance from** the perturbed **magnetopause** (at $\psi = 0$).
- $\Rightarrow \langle \Delta\psi^2 \rangle$ over all lines **measures** the **effective magnetic diffusion** due to **reconnection**.
- ▶ Effective $D_{\text{eff}} \approx 10^{10} \text{ m}^2/\text{s}$:)))

A minimal plasma model: resistive Hall-MHD

Continuity equation

$$\partial n / \partial t + \nabla \cdot (n \vec{U}) = 0$$

Momentum equation

$$\partial(n \vec{U}) / \partial t + \nabla \cdot [(n \vec{U} \vec{U}) + (P \overleftrightarrow{T} - \vec{B} \vec{B})] = 0 ; P = P_i + P_e + |\mathbf{B}|^2 / 2$$

Adiabatic closures

$$\partial(n S_{i,e}) / \partial t + \nabla \cdot (n S_{i,e} \vec{u}_{i,e}) = 0 ; S_{i,e} = P_{i,e} n^{-5/3}$$

Faraday & current equations

$$\partial \vec{B} / \partial t = -\nabla \times \vec{E} ; \vec{J} = \nabla \times \vec{B}$$

Generalized Ohm's law

$$\vec{E} = \underbrace{-\vec{U} \times \vec{B} + \vec{J} / n \times \vec{B}}_{-\vec{U}_e \times \vec{B}} - \frac{1}{n} \nabla P_e + \eta \vec{J}$$

Boundary conditions

► **MHD characteristic decomposition** at the x -boundaries

$$(L_a^\pm, L_s^\pm, L_f^\pm, L_0) \leftrightarrow (\rho, T, v, B_y, B_z)$$

$$\partial/\partial t (\rho, T, v, B_y, B_z) = F(L_a^\pm, L_s^\pm, L_f^\pm, L_0)$$

$$(L_a^\pm, L_s^\pm, L_f^\pm, L_0) = G(a, s, f, \rho, T, v, B_y, B_z, \partial_x)$$

► **Non-reflective** boundary conditions (left boundary) and **equilibrium sustainment**:^a

$$\rightarrow L_{0,a,s,f}^\pm = L_{0,a,s,f}^\pm|_{\text{internal points}} \text{ for outgoing waves}$$

$$\leftarrow L_{0,a,s,f}^\pm = L_{0,a,s,f}^\pm|_{\text{equilibrium}} \text{ for incoming waves}$$

^aFaganello *NJP* 2009

Late nonlinear evolution

